

(37-38)₇ Monday [2014-04-07]

16:00 - 18:00 413

Material response

$$T = \hat{T}(F) \quad \text{elastic material}$$

↑ response function

$$p_A = p_0 + F(\bar{p}_A - \bar{p}_0)$$

$$p_A^* = q_0^* + Q(p_A - p_0) \quad \text{change of observer}$$

$$p_0^* = q_0^* + Q(p_0 - p_0) = q_0^*$$

$$p_A^* - p_0^* = Q(p_A - p_0) = QF(\bar{p}_A - \bar{p}_0) = F^*(\bar{p}_A - \bar{p}_0)$$

$$T^* = QTQ^T$$

$$T^* = \hat{T}(F^*)$$

$$F^* = QF$$

$$\hat{T}(F^*) = Q\hat{T}(F)Q^T \quad \forall Q$$

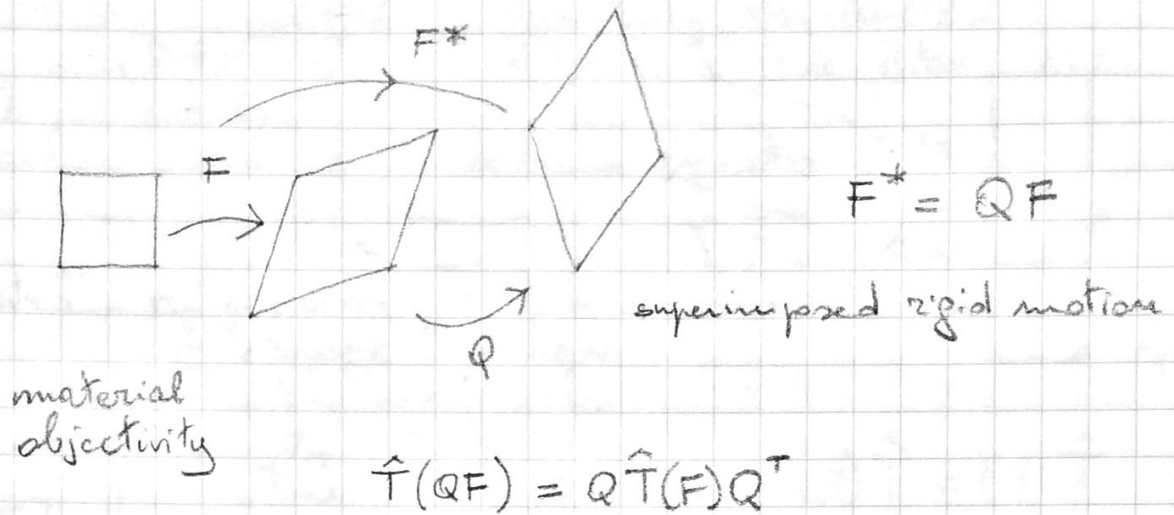
$$\hat{T}(QF) = Q\hat{T}(F)Q^T \quad \forall Q, \forall F$$

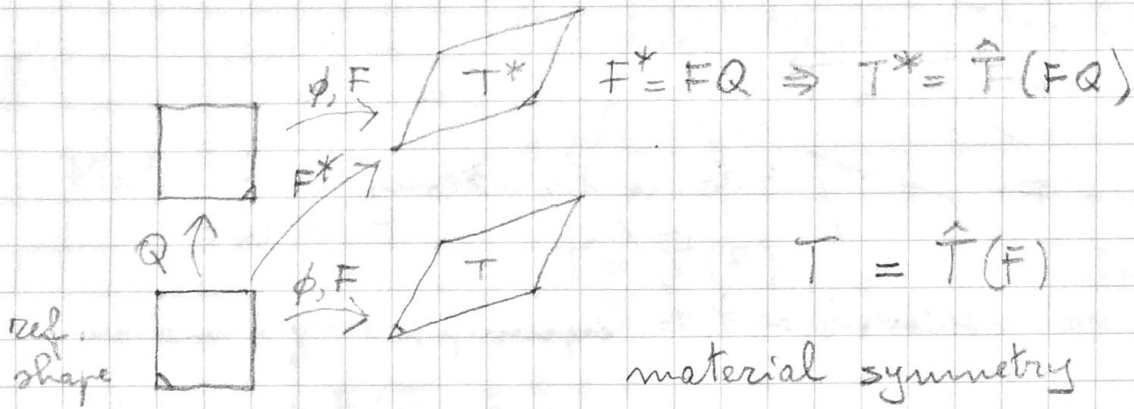
$$\hat{T}(QRU) = Q\hat{T}(RU)Q^T \quad \forall Q$$

$$Q = R^T \Rightarrow \hat{T}(U) = R^T\hat{T}(F)R$$

$$\hat{T}(F) = R\hat{T}(U)R^T$$

objective response function = reduced form





$$\hat{T}(F) = \hat{T}(FQ) \quad \hat{T} \text{ is } Q \text{ invariant}$$

Symmetry group: collection of such rotations Q

Isotropic material: symmetry group $O(3)^+$

Isotropy

$$\hat{T}(F) \stackrel{\text{objectivity}}{=} \hat{T}(FQ) = \hat{T}(RUQ) = \hat{T}(\underbrace{R}_{\tilde{R}} \underbrace{Q}_{\tilde{U}} \underbrace{Q^T}_{\tilde{U}} \underbrace{U}_{\tilde{U}} Q) \quad \forall Q$$

$$= (RQ) \hat{T}(Q^T U Q) (RQ)^T = R \hat{T}(U) R^T$$

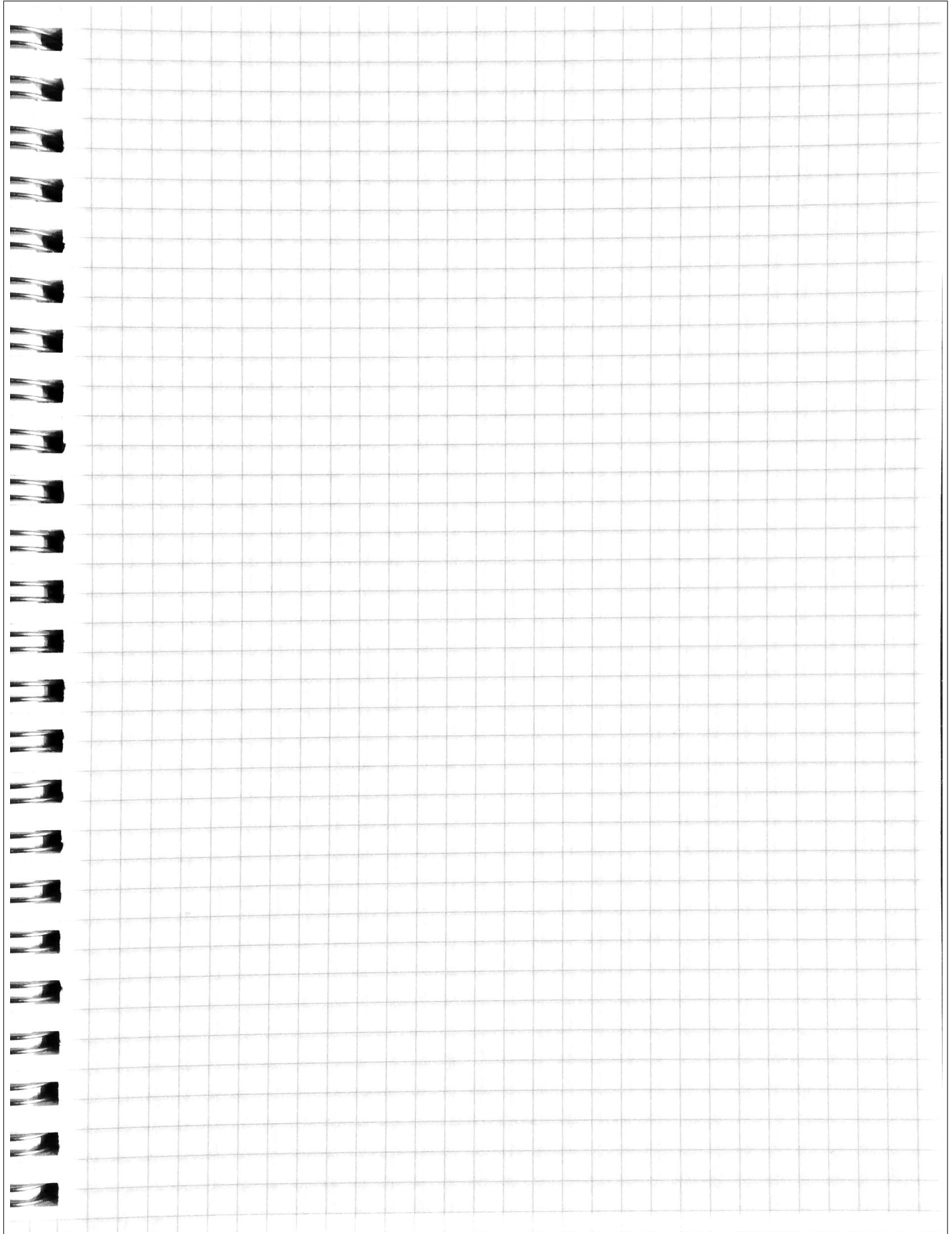
$$= R Q \hat{T}(Q^T U Q) Q^T R^T$$

$$\hat{T}(F) = \hat{T}(RU) = R \hat{T}(U) R^T$$

$$\Rightarrow Q \hat{T}(Q^T U Q) Q^T = \hat{T}(U)$$

$$\Leftrightarrow \hat{T}(Q^T U Q) = Q^T \hat{T}(U) Q \quad \forall Q$$

Isotropic response function



(39-40)₇ Tuesday [2014-04-08]
9:00-11:00

$$T \cdot L V_R = T \cdot L V_R \det F$$

$$L (P_A - P_0) = LF (\bar{P}_A - \bar{P}_0)$$

$$\begin{aligned} T \cdot L V_R \det F &= S \cdot LF V_R \\ &= SF^T \cdot L V_R \end{aligned}$$

$$T \det F = SF^T$$

$$S = TF^{-T} \det F$$

PIOLA stress

$$\int_{\mathcal{R}} T \cdot \nabla_{\mathbf{r}} dV = \int_{\bar{\mathcal{R}}} (T \cdot \nabla_{\mathbf{r}}) \det F dV = \int_{\bar{\mathcal{R}}} S \cdot \nabla_{\bar{\mathbf{r}}} dV$$

$$\nabla_{\bar{\mathbf{r}}} = \nabla_{\mathbf{r}} F \quad \Rightarrow \quad T \cdot \nabla_{\mathbf{r}} \det F = T \cdot \nabla_{\bar{\mathbf{r}}} F^{-T} \det F$$

$$= (\det F) (TF^{-T}) \cdot \nabla_{\bar{\mathbf{r}}} = S \cdot \nabla_{\bar{\mathbf{r}}}$$

$$\operatorname{div}(S^T \bar{r}) = S \cdot \nabla \bar{r} + \operatorname{div} S \cdot \bar{r}$$

fields on $\bar{\mathcal{R}}$

$$\int_{\mathcal{R}} b \cdot r \, dV = \int_{\bar{\mathcal{R}}} (b \cdot r) \det F \, dV = \int_{\bar{\mathcal{R}}} \bar{b} \cdot \bar{r} \, dV$$

$$\bar{r}(x) = r(\phi(x)), \quad \bar{b} = b \det F$$

$$\int_{\partial \mathcal{R}} t \cdot r \, dA = \int_{\partial \bar{\mathcal{R}}} \bar{t} \cdot \bar{r} \, dA \quad \bar{t} = t \frac{A_{\bar{\mathcal{R}}}}{A_{\mathcal{R}}}$$

$$\int_{\mathcal{R}} T \cdot \nabla r \, dV = \int_{\bar{\mathcal{R}}} S \cdot \nabla \bar{r} \, dV$$

$$\int_{\bar{\mathcal{R}}} \bar{b} \cdot \bar{r} \, dV + \int_{\partial \bar{\mathcal{R}}} \bar{t} \cdot \bar{r} \, dV - \int_{\bar{\mathcal{R}}} S \cdot \nabla \bar{r} \, dV = 0 \quad \forall \bar{r}$$

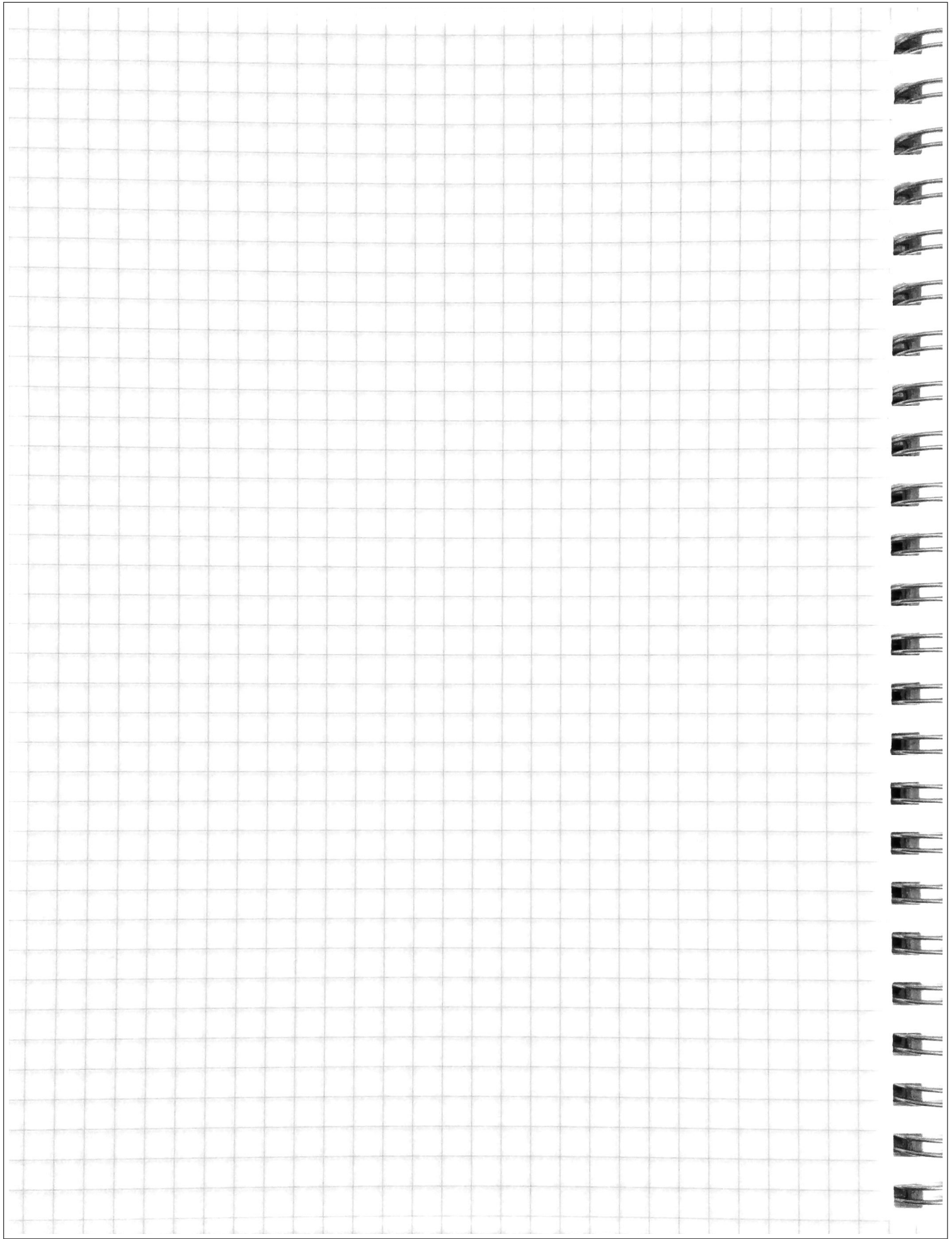
$$[\dots] \int_{\bar{\mathcal{R}}} \operatorname{div}(S^T \bar{r}) \, dV = \int_{\partial \bar{\mathcal{R}}} S^T \bar{r} \cdot \bar{n} \, dV$$

↑ outward unit normal

$$[\dots] \begin{cases} \bar{b} + \operatorname{div} S = 0 & \bar{\mathcal{R}} \\ \bar{t} - S \bar{n} = 0 & \partial \bar{\mathcal{R}} \end{cases}$$

compare

$$t = T n, \quad \bar{t} = S \bar{n} \quad S = T \operatorname{cof} F$$



(41-42)₇ Wednesday [2014-04-09]

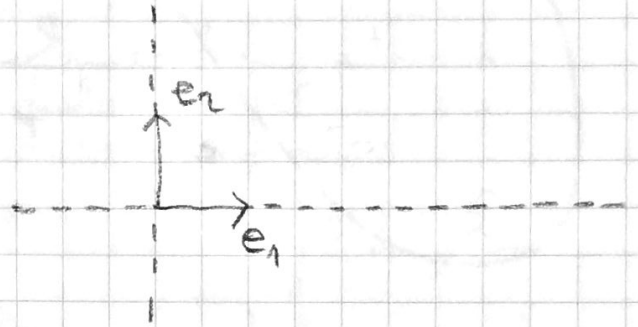
8:00 - 11:00

A1,2

Scalar form of the balance equations

$$\operatorname{div} T \cdot e_i = \operatorname{div}(T^T e_i) = \operatorname{tr} \nabla(T^T e_i)$$

$$[T] = \begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{pmatrix}$$



$$T^T e_1 = \sigma_{11} e_1 + \sigma_{12} e_2 + \sigma_{13} e_3$$

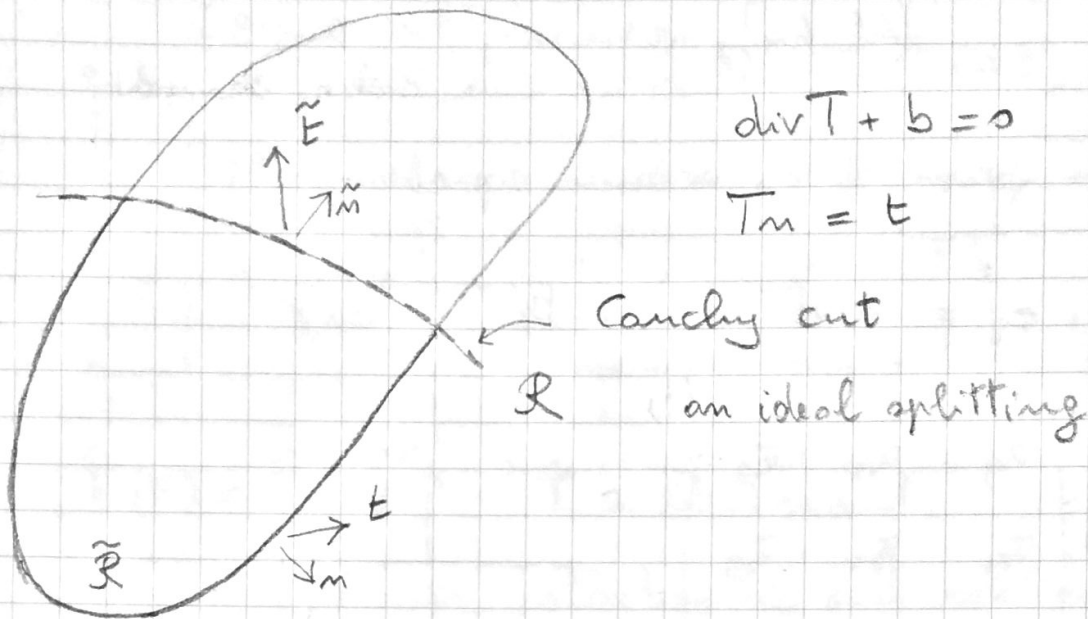
$$T^T e_2 = \sigma_{21} e_1 + \sigma_{22} e_2 + \sigma_{23} e_3$$

$$T^T e_3 = \sigma_{31} e_1 + \sigma_{32} e_2 + \sigma_{33} e_3$$

$$[\nabla(T^T e_i)] = \begin{pmatrix} \sigma_{11,1} & \sigma_{11,2} & \sigma_{11,3} \\ \sigma_{12,1} & \sigma_{12,2} & \sigma_{12,3} \\ \sigma_{13,1} & \sigma_{13,2} & \sigma_{13,3} \end{pmatrix}$$

$$\operatorname{div} T \cdot e_1 = \sigma_{11,1} + \sigma_{12,2} + \sigma_{13,3}$$

$$\operatorname{div} T \cdot e_i = \sigma_{i1,1} + \sigma_{i2,2} + \sigma_{i3,3}$$

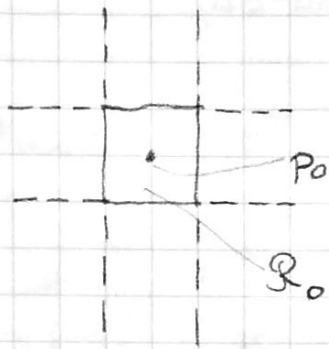


$$\int_{\tilde{R}} b \cdot r \, dV + \int_{\partial \tilde{R}} t \cdot r \, dA - \int_{\tilde{R}} T \cdot \nabla r \, dV = 0$$

$$[\dots] \quad \operatorname{div} T + b = 0$$

$$Tm = t, \quad T\tilde{m} = \tilde{t}$$

\tilde{t} is the traction we should apply to keep the lower part of the body balanced after removing the upper part.



we can cut a sample
in the shape of a cube centered
at p_0 : $R_0 \subset R$

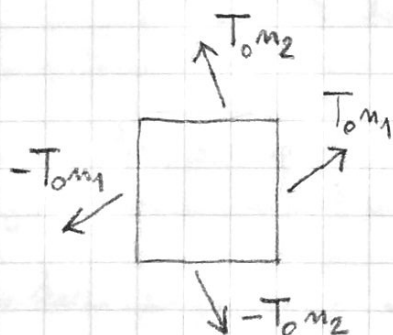
By using an affine test velocity field we get

$$\int_{R_0} T \cdot \nabla v \, dV = \int_{R_0} T \cdot L \, dV = \left(\int_{R_0} T \, dV \right) \cdot L$$

If we compare this expression with the one we
wrote for an affine body we find

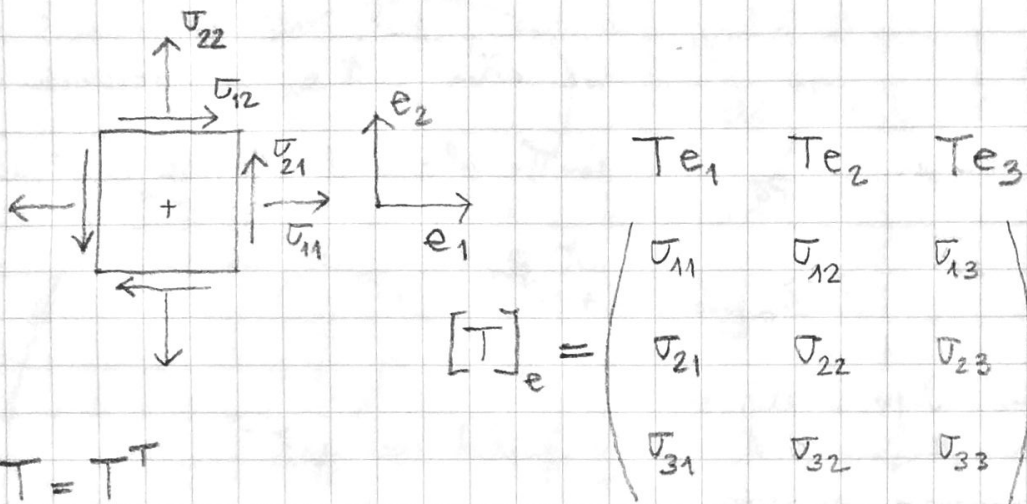
stress in the
affine body $T_0 \cdot L \cdot V_{R_0} = \int_{R_0} T \, dV \cdot L$
stress field

$$\Rightarrow T_0 = \frac{1}{V_{R_0}} \int_{R_0} T \, dV \quad \text{mean stress}$$



traction on the boundary of the cube
balanced by the uniform stress

$$T_0 \approx T(p_0) \quad (\text{small } V_{R_0})$$



$$T = T^T$$

$$\Rightarrow \sigma_{21} = \sigma_{12} \text{ [!!!]}$$

↓
real eigenvalues and orthogonal eigenvectors

$$\sigma_1, \sigma_2, \sigma_3$$

principal stresses

$$a_1, a_2, a_3$$

principal directions

diagonal matrix $[T]_a = \begin{pmatrix} \sigma_1 & & \\ & \sigma_2 & \\ & & \sigma_3 \end{pmatrix}$

